Integrable systems in perturbative quantum field theory

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ENIGMA 08

Rutger Boels (NBIA)

Integrable systems in perturbative QFT

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A (1) > A (2) > A

Largely based on

- 0805.1197 [hep-th] (with Christian Schwinn)
- hep-th/0702035, hep-th/0604040 (with David Skinner and Lionel Mason)
- 'the literature'

known

The self-dual Yang-Mills equations in 4 dimensions are integrable [Penrose and Ward]

main idea in our work

Full equations as a perturbation around self-dual sector [Mason, 05]

suspicion

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The big picture



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Integrable systems in perturbative QF1

- Standard model of particle physics is a Yang-Mills theory
- ullet ightarrow Yang-Mills equations predict experimental outcomes

Pure Yang-Mills theory

ingredients:

- A the gauge field: a Lie-algebra valued connection 1-form of a fibre bundle over ℝ⁴ (say U(N))
- $F = dA + g(A \wedge A)$ its curvature 2 form
- g the coupling constant
- $S = \operatorname{Tr} \int F \wedge *F$ the Yang-Mills action
- $\frac{\partial S}{\partial A} = 0$ are the Yang-Mills equations
- equations non-linear if $g \neq 0$
- boundary conditions describe the experiment/process under study

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- want: solve Yang-Mills equations with plane-wave asymptotics.
- \rightarrow scattering amplitudes in perturbative series in g
- well-defined procedure to calculate these in terms of path integrals (cf. math-ph/0204014)
- scattering data: Lorentz group rep (P_{μ} and polarization)
 - Hodge * projects 2 forms onto ± 1 eigenvalues: self-dual and anti-selfdual.
 - (anti) self-dual connections A solve full Yang-Mills equations.
 - example: polarization of external fields (helicity)

surprise

some scattering amplitudes are very, very simple [80's]

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$$A(++\ldots+) = 0, A(-+\ldots+) = 0, A(--+\ldots+) =$$
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• despite rapid growth ($\sim n^2$) of ordinary diagram complexity

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selfdual solutions to the Yang-Mills equations on \mathbb{R}^4 (S⁴) (on to one) holomorphic flat vector bundles on \mathbb{CP}^3

• CP³ is the projective twistor of four dimensional space

- Euclidean: locally, $CP^3 = R^4 \times \mathbb{CP}^1$
- (see [Woodhouse, 85]) ingredients: (0, 1) forms, Lie algebra valued connection a, weight 0, adjoint valued b, weight -4

$$S = \operatorname{Tr} \int_{\mathbb{CP}^3} d\Omega \wedge b \wedge \left(\overline{\partial} a + ga \wedge a\right)$$

- *b*-field equation: F[a] = 0 (zero curvature)
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• 4*D* field theory (known as self-dual YM) as an integrable system puzzle:

• Space-time SD YM \rightarrow one non-trivial scattering amplitude A(++-) (classically), while the twistor system is free

resolution:

• Scattering amplitude is part of the field transformation (cf. KdV) [Ettle-Fu-Fudger-Mansfield-Morris, 07]

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summary of knowledge up to 2003

$$\begin{array}{c} \mathbb{CP}^3 \quad \int_{\mathbb{CP}^3} d\Omega \wedge b \wedge \left(\overline{\partial} a + g a \wedge a \right) & + ? \\ \mathbb{R}^4 \quad \int_{\mathbb{R}^4} B \wedge F_+ & \int_{\mathbb{R}^4} F \wedge *F \end{array}$$

- In dec. 2003 Witten proposed 'twistor string theory' (also [Nair, 88])
- Amplitudes from topological B-model on supertwistor space CP^{3|4} ('super Calabi-Yau')
- Very inspiring in the physical community (~ 400 citations)
- [Mason 2005] and our work provides missing term:

 $S_{\text{miss}} = \int_{\mathbb{R}^{4} \times \mathbb{CP}^{1} \times \mathbb{CP}^{1}} d^{4}x \wedge dk_{1} \wedge dk_{2} \wedge (HbH^{-1})_{1} \wedge (HbH^{-1})_{2} \langle 12 \rangle^{2}$

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$$(\overline{\partial}_{0} + a_{0})H[a_{0}] = 0$$

$$S_{\text{miss}} \sim \sum_{i=0}^{\infty} \left(\sum_{j=1}^{i} ba \dots b_j \dots a \right)$$

- use same transformation to action angle coordinates as before
- interpolates between usual Yang-Mills eqns and 'CSW rules'
- generates those 'simple' scattering amplitudes mentioned before
- re-orders complicated perturbation theory
- \rightarrow perturbation around an integrable system!

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$$S_{\text{miss}} = \int_{\mathbb{R}^{4} \times \mathbb{CP}^{1} \times \mathbb{CP}^{1}} d^{4}x \wedge dk_{1} \wedge dk_{2} \wedge (HbH^{-1})_{1} \wedge (HbH^{-1})_{1} \langle 12 \rangle^{2}$$

$$(\overline{\partial}_0 + a_0)H[a_0] = 0$$

 $S_{\mathrm{miss}} \sim \sum_{i=0}^{\infty} \left(\sum_{j=1}^i ba \dots b_j \dots a\right)$

- use same transformation to action angle coordinates as before
- interpolates between usual Yang-Mills eqns and 'CSW rules'
- generates those 'simple' scattering amplitudes mentioned before
- re-orders complicated perturbation theory
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- twistor variables provide a transform to action-angle variables \rightarrow trivializes 1 three particle amplitude
- Yang-Mills has two three particle amplitudes (*A*(+ + -) and *A*(+ -))
- 'Other' three particle amplitude simplified on 'dual' twistor space
- Can you simplify both at the same time?

→ ambi-twistor space? [Mason, Skinner, 05]

BCFW recursion

Britto-Cachazo-Feng-Witten (05): every amplitude in Yang-Mills can be expressed in terms of certain sums over those 2 three particle amplitudes only.

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- Is full Yang-Mills integrable?
- BCFW: reconstruction of scattering amplitudes through a field transformation?

More interesting ideas floating around

- similar constructions for physically interesting cases
 - supersymmetry (very natural)
 - massive gauge theories
- beginning of understanding of the quantum aspects (much harder)
- similar construction for Einstein gravity [Mason, Skinner, 08]

BCFW remarks

- BCFW works in 10 dimensions [Arkhani-Hamed, Kaplan, 08]
- recursive BCFW structure is very natural in string theory [Boels, Larsen, Obers, Vonk, 08]
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Interesting developments from/with physical motivationNew input/life for twistor theory and ideas

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- Is full Yang-Mills classically integrable?
- Is full (supersymmetric?) Yang-Mills quantum mechanically integrable?
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